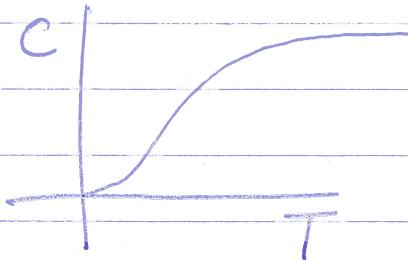


3/42

Fri. Aug. 31

(1)

- We have been discussing the "Einstein Solid" which consists of independent quantum harmonic oscillators
- Classical gives Dulong-Petit law
- QHO's explain non-linear behavior of heat capacity at low temperatures
- I'd like to look at one more parameter in the fit for C
- all oscillators have one frequency ω , sometimes called "Einstein" frequency
- in the C vs. T curve: $C(T) = 3k_B (\beta \hbar \omega)^2 \frac{e^{\beta \hbar \omega}}{(e^{\beta \hbar \omega} - 1)^2}$



- $T_E = \hbar \omega / k_B$ is a fit parameter that gives the best fit to the exp. data

$T_E \propto \omega$ or ω_E , both constants

(2)

- Example being for copper (Cu)

$$T_E = 240\text{K} \quad (-33^\circ\text{C}) \quad * \text{ record low Flgstaff}$$

-34°C

- above T_E , we get more-or-less Dulong-Petit law

- one well known discrepancy with the classical law $C/R \approx 3$, most materials for diamond

$$\frac{C}{R} \approx 0.7$$

• heat capacity per mole

at room T. Why?

so room T
is "cold"

$$T_{E, \text{diamond}} = 1320\text{K} > \text{room T}$$

- why does it have such a value?

- due to the chemical bonds between C atoms

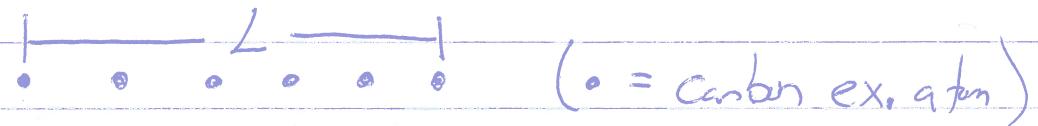
- the bond is strong (think of hardness of diamond) + carbon atoms are light

$$\omega = \sqrt{\frac{K}{m}}, \quad K \text{ large } \left\{ \begin{array}{l} \omega \text{ large} \\ m \text{ small} \end{array} \right\} \Rightarrow T_E \text{ large}$$

(3)

- This discussion of bonds between atoms leads us directly to the idea of coupled harmonic oscillators

- let's look at a 1D "solid"



- transverse modes \uparrow
- long. modes \leftrightarrow
- we usually use "periodic boundary conditions"

- Circumference L
- \dots
- \dots
- \dots
- \dots

- first atom is coupled to last, etc.

- translate along solid by L distance, we end up at the same place

$$\text{• a wave } e^{ikr} = e^{ik(r+L)} \quad (\text{wave is same at } r + r + L)$$

Note: • We are looking at continuous medium with large L (forget atoms for now!) \rightarrow

(4)

$$\Rightarrow e^{ikr} = e^{ikr} e^{ikL}$$

$$e^{ikL} = 1 = \cos(kL) + i\sin(kL)$$

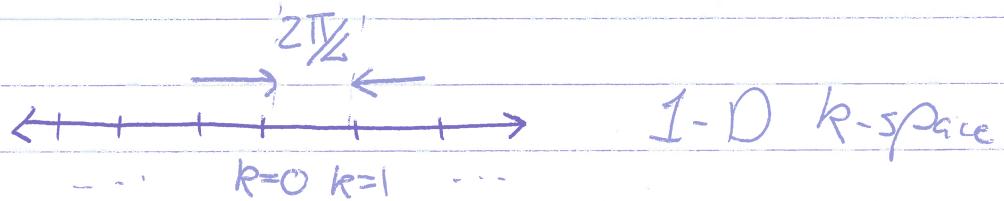
- this is true only when $kL = 2\pi n$
- in other words only modes with wave number

$$k = \frac{2\pi n}{L}$$

are possible (recall $k = 2\pi/\lambda$)

• can be thought of ^{as} a "spatial frequency"

• in solid state, we use "k-space" a lot



• each point on axis is a mode (vibration)

• for large L, $\Delta k = 2\pi/L$ is small

• for solids, L is large (many atoms) ^{end}
g-31