

17/42

Fri. Oct. 05

①

- We have been investigating the dispersion of a 1D solid with discrete nature (atoms)
- Using periodic boundary conditions, we see that

$$X_{n+N} = X_n$$

where N is the # of atoms &
 X_n is the position of atom 'n'

- we see that $U_n = U_{n+N}$ on $\xrightarrow{\text{disp. from eq.}}$

$$Ae^{i(kX_n - \omega t)} = Ae^{i(kX_{n+N} - \omega t)}$$

- recall $X_n = na$ & $X_{n+N} = (N+n)a$

$$\Rightarrow Ae^{i(kna - \omega t)} = Ae^{i(k[N+n]a - \omega t)}$$

$$\Rightarrow e^{ikNa} = 1$$

$$\Rightarrow \cos(kNa) + i \sin(kNa) = 1$$

$$kNa = 2\pi p \quad (p \in \mathbb{Z} \text{ integers})$$

(2)

- therefore $k = \frac{2\pi p}{Na}$ or $k = \frac{2\pi p}{L}$
- where $L = Na$ is the total length of the solid
- so again, k values are discrete with spacing between k 's being

$$\frac{2\pi}{L}$$

- we only count modes in first Brillouin zone

$$-\pi/a \leq k < \pi/a \quad (\text{note don't double count mode at max. } \omega)$$

$$N = \frac{\text{range of } k\text{'s}}{\text{spacing between } k\text{'s}} \quad (\# \text{ of modes})$$

$$= \frac{2\pi/a}{2\pi/Na} = N$$

- just as was predicted by Debye //

③

- the energy of a lattice vibration is quantized & a quantum of such energy is called a phonon.
- given a mode with frequency ω , the energy is given by

$$E = \left(\frac{1}{2} + n\right) \hbar \omega$$

for excited state 'n', on the mode is occupied by 'n' phonons (quanta), $E = \hbar \omega$

- in analogy to classical physics, higher energy corresponds to larger amplitude of vibration (louder)
- each step up the ladder is a phonon
- we may have any number of phonons in the same state (bosons), just like photons
- Bose occupation factor

$$n_B(\beta \hbar \omega) = \frac{1}{e^{\beta \hbar \omega} - 1}$$

④

- so the expected value for a given mode

$$E_k = \hbar \omega(k) \left[n_B(\beta \hbar \omega(k)) + \frac{1}{2} \right]$$

- and of course the total energy is the sum of energies E_k

$$E_{\text{total}} = \sum_k \hbar \omega(k) \left[n_B(\beta \hbar \omega(k)) + \frac{1}{2} \right]$$

- the sum over k is for the first Brillouin zone only

- we know for real solid N is large

$$\sum_k \rightarrow \frac{Na}{2\pi} \int_{-\pi/a}^{\pi/a} dk$$

- check: $\frac{Na}{2\pi} \int_{-\pi/a}^{\pi/a} dk = N$

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- thus the total energy,

$$E_{\text{total}} = \frac{N_a}{2\pi} \int_{-\pi/a}^{\pi/a} dk \hbar \omega(k) \left[n_B(\beta \hbar \omega(k)) + \frac{1}{2} \right]$$

- From which we may obtain the heat capacity

* note
See
page 5.5

- extending the analogy from photons, phonons also have momentum

$$\vec{p} = \hbar \vec{k}$$

which is called the "crystal momentum"

- this formulation allows us to treat phonons as particles
- For example, we see that phonons scatter from each other or phonons/photons interact etc.

- recall that for the continuum model
the density of states

$$g(\omega) = \frac{L}{\pi} \frac{1}{d\omega/dk}$$

- for the 1D lattice:

$$g(\omega) = \frac{2L}{\pi a \omega_{\max}} \left[\cos\left(\frac{ka}{2}\right) \right]^{-1}$$

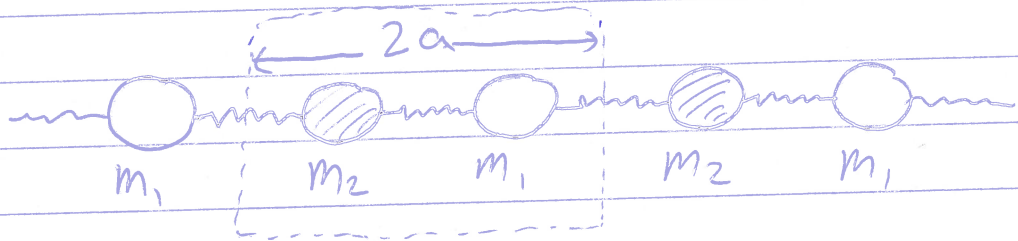
- Heat capacity

$$C_v = k_B \int \left(\frac{\hbar\omega}{k_B T} \right)^2 e^{\hbar\omega/k_B T} \left(e^{\hbar\omega/k_B T} - 1 \right)^{-2} g(\omega) d\omega$$

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- although it might seem pedantic to worry about what happens for a diatomic molecule, but some important physics is there

- Model: two repeating atoms of mass m_1 & m_2 in repeating units



- spring constant α is same
- the dashed box is our first encounter with a "unit cell" here of length $2a$ (periodicity)
- Similar to our monatomic model we now have 2 diff. eqs.

Over